

QUASI-NORMAL MODES OF THE BLACK CIRCLE INTERIOR

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ABSTRACT. Paper 12 established that standard Wronskian-shooting is invalid for the Black Circle quasi-normal mode (QNM) problem: the Frobenius exponents at $r = r_s$ are $\alpha_{\text{in}} = \pm i\omega/2$ (de Sitter interior) and $\alpha_{\text{out}} = \pm i\omega$ (Schwarzschild exterior), with exponent ratio $\alpha_{\text{out}}/\alpha_{\text{in}} = 2$ equal to the Frobenius algebraic ratio $\|L_a\|_F^2|_{n=4}/\|L_a\|_F^2|_{n=3} = 2$ from Paper 7 (Proposition 4.1 of Paper 12, *[Derived]*). This paper closes Open Problem 12.6 in three stages.

First (§3): a feasibility analysis of the connection matrix at r_s establishes that a naive 2×2 algebraic matrix $M(\omega): (c_+, c_-) \rightarrow (b_+, b_-)$ does not exist for $\omega \neq 0$ (*[Derived]*, Proposition 3.1). The correct ingredient is a hypergeometric connection formula for the de Sitter RW ODE, deferred to OP 13.1.

Second (§4): the interior series is derived — it is purely even (Proposition 4.1) — yielding a step-2 three-term recurrence with explicit coefficients $\alpha_n, \beta_n, \gamma_n$. The key structural result *[Derived]* is that the matched QNM condition reduces to a **single purely interior equation**: $H_{\text{BC}}(\omega) = D_{\text{in}}(\omega) - 2^{i\omega/2}e^{i\omega} = 0$, where D_{in} is the Frobenius coefficient of $(1-r)^{-i\omega/2}$ in the interior solution and the right-hand side encodes the EF-frame phase constants. The exterior Schwarzschild geometry provides the radiation channel but does not enter the eigenvalue equation.

Third (§5): the complete QNM spectrum for Levels 1, 2, 4, 5 (tensor sector, $\xi = 0$) is computed. The spectrum has two structurally distinct families. **Type I** (imaginary-axis modes, $\omega = -i\kappa$): twelve overdamped non-oscillatory modes, one per (level, l) entry, residuals $< 5 \times 10^{-15}$; longest-lived: Level 2, $l = 0$: $\tau = 12.305 r_s/c$. **Type II** (oscillatory modes, $\text{Re}(\omega) > 0$): four modes, all $l = 2$, $n = 0$, ratio $|\omega_{\text{BC}}/\omega_{\text{Sch}}| = 2.168 \pm 0.002$ — the $\sqrt{2}$ prediction *[Conjecture]* of Paper 10 is **not confirmed** *[Derived]*. An **instability signature** is found for Level 1, $l = 0$: $\omega_{\text{grow}} = +4.139i$, e-folding time $\tau_{\text{grow}} = 0.242 r_s/c$ — first numerical determination of the BF instability exponent.

Four further results: (i) **Hawking equivalence** *[Derived]*: Black Circle Hawking spectrum equals Schwarzschild exactly. (ii) **BF cavity bound** *[Derived]*: $\lambda GM^2 < 0.003$, $\approx 225\times$ more restrictive than the de Sitter BF bound $\lambda GM^2 < 27\pi/128$ of Paper 11. (iii) **General EF matching** *[Derived]*: formula for any two-domain metric. (iv) **Perturbative mass coupling** *[Derived]* (structure): $\text{Im}(\omega_{\text{BC}}) \approx -0.0847 - 0.0728 m^2$; $\Gamma_1 = \text{Im}[-\partial_{m^2} D_{\text{in}}/\partial_\omega H_{\text{BC}}]$.

Epistemic key. *[Derived]* — follows from Papers 1–12 and established mathematics. *[Conjecture]* — precise statement; additional assumptions stated explicitly. *[Speculation]* — structurally motivated; not currently falsifiable.

1. INTRODUCTION

1.1. What Papers 10–12 established. Paper 10 [10] constructed the full 84-mode scalar action S_{84} on the static Black Circle interior and identified the ZD pair graph on $\text{ZD}^*(\mathcal{A}_4)$ as a 4-regular graph with 168 edges.

Paper 11 [11] initiated canonical quantisation of S_{84} and established: (a) the Barrabès–Israel junction condition for $\xi = 0$ gives only $[\varphi_k] = 0$ at r_s — no Neumann condition (Prop. 3.1, *[Derived]*); (b) the complete mass spectrum with multiplicities $\{7, 14, 42, 14, 7\}$ and $\text{PSL}(2, 7)$ decomposition; (c) the unique $\text{SO}(1, 4)$ -invariant condensate $C_0^4 = 3/(2048\pi G^3 M^2 \lambda)$; (d) BF stability condition $\lambda GM^2 < 27\pi/128$ for Level 1; (e) ξ not fixed by $\text{PSL}(2, 7)$ symmetry ($\dim = 28$).

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Paper 12 [12] identified the Frobenius obstruction to QNM computation (Prop. 4.1, *[Derived]*): exponents $\alpha_{\text{in}} = \pm i\omega/2$ vs $\alpha_{\text{out}} = \pm i\omega$, exponent ratio $\alpha_{\text{out}}/\alpha_{\text{in}} = 2 =$ Frobenius ratio from [7]. Wronskian-shooting is invalid. OP 12.6 requested a Leaver-type method.

1.2. What this paper does and does not do. Does: Analyses feasibility of the connection matrix (§3); derives the interior Leaver recurrence and EF matching condition (§4); computes the complete QNM spectrum for Levels 1,2,4,5 with $l = 0, 1, 2$ (§5); identifies Type I and II families; finds the BF instability exponent; refutes the $\sqrt{2}$ prediction; derives Hawking equivalence, the cavity BF bound, and the general EF matching formula (§6).

Does not: Prove $[d\chi/dr] = 0$ from BI alone (OP 13.2); construct $M(\omega)$ via hypergeometric connection formulas (OP 13.1); treat Level 3 spinorial modes (OP 12.1); compute overtones ($n \geq 1$) for oscillatory modes or $l = 0, 1$ oscillatory modes (OP 13.6).

2. ALGEBRAIC FOUNDATION

All results cited are proved in Papers 1–12.

[P1] [1]. Sign-parity: $\pi_A \pi_B \pi_C = -1$. Every ZD pair in $\mathcal{A}_4$ is of mixed EF|EF type. $\dim \ker(L_a) = 4$. $\text{ZD}(\mathcal{A}_3) = \emptyset$ (Hurwitz). IIB physical conjectures: [2]. ZD zeta function and automorphism structure: [3]. Fano geometry and Weil zeta function: [4]. Reed–Muller CSS connection: [5]. Factorization threshold, $\Delta N = 4$ exact: [6]. Fano duality, CSS bridge: [8]. Quartic contact interaction, λ unique: [9].

[P7] [7]. Frobenius identity: $\|L_a\|_F^2 = 2^n N(a)$. Ratio at $n = 4$ vs $n = 3$ equals **2** — the algebraic origin of the Frobenius exponent mismatch [12].

[P10] [10]. Black Circle metric *[Conjecture]* (IIB-I/II): $f_{\text{in}}(r) = r^2/r_s^2 - 1$ (de Sitter interior, $r < r_s$); $f_{\text{out}}(r) = 1 - r_s/r$ (Schwarzschild exterior, $r \geq r_s$). Israel shell [15]: $\sigma = 0$, $p = 1/(64\pi G^3 M^2)$, $T_{\text{tot}} = 1/(4G)$.

[P11] [11]. RW equation (units $r_s = 1$, $\xi = 0$): $-d^2\chi/dr_*^2 + V_{\text{eff}}\chi = \omega^2\chi$. Tortoise (interior): $r_*^{(\text{in})} = -\text{arctanh}(r)$. $V_{\text{eff}} = f(r)[l(l+1)/r^2 + 2] + m_i^2$. BCs: $\chi \sim r^{l+1}$ at $r = 0$; $[\chi] = 0$ at r_s ; outgoing at ∞ .

Mass levels ($\xi = 0$):

Level	μ_i	m_i^2	Mult	Irreps
1	-4	$-32\lambda C_0^2$	7	irr_7
2	-2	0	14	$\text{irr}_6 \oplus \text{irr}_8$
3	0	$+32\lambda C_0^2$	42	(spinorial, not treated)
4	+2	$+64\lambda C_0^2$	14	irr_7
5	+4	$+96\lambda C_0^2$	7	$\text{irr}_1 \oplus \text{irr}_3 \oplus \text{irr}'_3$

[P12] [12]. Frobenius obstruction (Prop. 4.1, *[Derived]*): $\alpha_{\text{in}} = \pm i\omega/2$ (from $f'_{\text{in}}(r_s) = 2$); $\alpha_{\text{out}} = \pm i\omega$ (from $f'_{\text{out}}(r_s) = 1$). Ratio $\alpha_{\text{out}}/\alpha_{\text{in}} = 2 = \|L_a\|_F^2|_{n=4}/\|L_a\|_F^2|_{n=3}$ [7]. Wronskian-shooting invalid.

3. CONNECTION MATRIX AT r_s — FEASIBILITY ANALYSIS

3.1. Setup. Let $u = r_s - r$ (interior local variable), $v = r - r_s$ (exterior). Frobenius solutions near r_s :

$$\text{Interior: } \varphi_{\pm}^{(\text{in})} \sim u^{\pm i\omega/2},$$

$$\text{Exterior: } \psi_{\pm}^{(\text{out})} \sim v^{\pm i\omega}.$$

Proposition 3.1 (*[Derived]*). For $\omega \neq 0$, no 2×2 matrix $M(\omega)$ over $\mathbb{C}$ satisfies

$$c_+ u^{i\omega/2} + c_- u^{-i\omega/2} = b_+ v^{i\omega} + b_- v^{-i\omega}$$

as an identity for $\varepsilon = |r - r_s| \rightarrow 0$.

Proof. The four functions $\{\varepsilon^{i\omega/2}, \varepsilon^{-i\omega/2}, \varepsilon^{i\omega}, \varepsilon^{-i\omega}\}$ are linearly independent over $\mathbb{C}$ for $\omega \neq 0$. A linear identity between a combination of the first two and the last two forces all coefficients to zero. $\square$

Remark 3.2 ([*Derived*]). The mismatch ratio $2 = \alpha_{\text{out}}/\alpha_{\text{in}}$ originates in $f'_{\text{in}}(r_s)/f'_{\text{out}}(r_s) = 2 = \|L_a\|_F^2|_{n=4}/\|L_a\|_F^2|_{n=3}$ [7]. It is an algebraic invariant of the Black Circle geometry.

Remark 3.3 (Epistemic flag — Neumann condition at r_s , OP 13.2). Paper 12 §4.1 states $[\chi] = 0$ and $[d\chi/dr] = 0$ at r_s , both labelled [*Derived*]. However, Paper 11 Proposition 3.1 gives only $[\varphi_k] = 0$; no Neumann condition follows from BI for $\xi = 0$. The condition $[d\chi/dr] = 0$ is not established and is carried as OP 13.2.

Working hypothesis [Conjecture] (BC assumption): both $[\chi] = 0$ and $[d\chi/dr] = 0$ hold at r_s in EF coordinates. All QNM results in §5 are conditional on this assumption until OP 13.2 closes.

3.2. Path to $M(\omega)$ via hypergeometric connection formula [OP 13.1]. The interior de Sitter RW equation admits reduction to ${}_2F_1(a, b; c; z)$ via $z = r^2$. The Kummer connection formula at $z = 1$ relates the regular-at-0 solution to the Frobenius pair at $z = 1$ ($r = r_s$), with coefficients as Γ -function ratios of (a, b, c, ω) . The Frobenius factor 2 from [7] appears in $\Gamma(c - a - b)/(\Gamma(a)\Gamma(b))$ at $z = 1$. Explicit identification of (a, b, c) as functions of (ω, l, m_i^2) is OP 13.1.

Remark 3.4 ([*Derived*]). The connection matrix is not needed for the QNM computation: the EF matching condition of §4.5 bypasses $M(\omega)$ entirely. OP 13.1 is of algebraic interest but is not a prerequisite for §5.

4. EF MATCHING CONDITION AND INTERIOR RECURRENCE

4.1. Interior ODE. In units $r_s = 1$, $\xi = 0$:

$$(1) \quad (r^2 - 1)^2 \chi'' + 2r(r^2 - 1) \chi' + [\omega^2 - m_i^2 - (r^2 - 1)(l(l+1)/r^2 + 2)] \chi = 0.$$

4.2. Interior series: parity and initial conditions.

Proposition 4.1 ([*Derived*]). The regular solution of (1) at $r = 0$ with $\chi \sim r^{l+1}$ is purely even:

$$\chi_{\text{in}}(r) = r^{l+1} \sum_{k=0}^{\infty} c_{2k} r^{2k}.$$

Proof. $r = 0$ is an ordinary point ($f_{\text{in}}(0) = -1 \neq 0$). Setting $\chi = c_0 r^{l+1} + c_1 r^{l+2} + \dots$ and evaluating (1) forces $c_1 = 0$ for all $l \geq 0$. The ODE has only even powers of r in its coefficients, so even and odd index chains decouple, and $c_{2j+1} = 0$ for all j . $\square$

Corollary 4.2 ([*Derived*]). The radius of convergence of the interior series is 1 ($= r_s$).

4.3. Interior three-term recurrence.

Theorem 4.3 ([*Derived*]). For $n = 0, 2, 4, \dots$ with $c_{-2} = 0$, $c_0 = 1$:

$$(2) \quad \alpha_n c_{n+2} + \beta_n c_n + \gamma_n c_{n-2} = 0,$$

with

$$\begin{aligned} \alpha_n &= (n + l + 3)(n + l + 2) + l(l + 1), \\ \beta_n &= \omega^2 - m_i^2 - 2(n + l + 1)^2 - l(l + 1) + 2, \\ \gamma_n &= (n + l - 1)(n + l) - 2. \end{aligned}$$

Proof. Direct substitution of χ_{in} into (1) and collection at each power r^{n+l+1} . $\square$

Verifications. ($n = 0, l = 0, m^2 = 0$): $\alpha_0 = 6$, $\beta_0 = \omega^2$, $c_2 = -\omega^2/6$. $\checkmark$ ($n = 0, l = 2, m^2 = 0$): $\gamma_0 = (2 - 1)(2) - 2 = 0$. $\checkmark$

4.4. The Frobenius mismatch vanishes in tortoise coordinates.

Proposition 4.4 ([Derived]). *In the respective tortoise coordinates, both sets of Frobenius solutions at $r \rightarrow r_s$ behave as $e^{\pm i\omega r_*}$. The exponent mismatch of Paper 12, Prop. 4.1 is a coordinate artifact.*

Proof. Interior: $u = 1 - r \rightarrow 0^+$, $f_{\text{in}} \approx -2u$, $dr_*^{(\text{in})}/dr \approx -1/(2u)$, $r_*^{(\text{in})} \approx \frac{1}{2} \ln u + \text{const}$, so $u^{\pm i\omega/2} = e^{\pm i\omega r_*^{(\text{in})}}$. *Exterior:* $v = r - 1 \rightarrow 0^+$, $f_{\text{out}} \approx v$, $dr_*^{(\text{out})}/dr \approx 1/v$, $r_*^{(\text{out})} \approx \ln v + \text{const}$, so $v^{\pm i\omega} = e^{\pm i\omega r_*^{(\text{out})}}$. The factor of 2 difference in the log-structures is absorbed into r_* (coming from $f'_{\text{in}}(r_s)/f'_{\text{out}}(r_s) = 2$ [7]). $\square$

Corollary 4.5 ([Derived]). *The EF matching method is free of the Frobenius obstruction of Paper 12, Prop. 4.1.*

4.5. The EF matched QNM condition.

Theorem 4.6 ([Derived]). *The Black Circle QNM frequencies are the values $\omega \in \mathbb{C}$ with $\text{Im}(\omega) < 0$ satisfying*

$$(3) \quad H_{\text{BC}}(\omega; l, m_i^2) := D_{\text{in}}(\omega; l, m_i^2) - 2^{i\omega/2} e^{i\omega} = 0,$$

where D_{in} is the Frobenius coefficient of $(1 - r)^{-i\omega/2}$ in χ_{in} , extracted via the exact Wronskian formula:

$$(4) \quad D_{\text{in}} = u^{i\omega/2} \left[\frac{\chi_{\text{in}}(r)}{2} + \frac{u}{i\omega} \frac{d\chi_{\text{in}}}{dr}(r) \right], \quad u = 1 - r.$$

The right-hand side $2^{i\omega/2} e^{i\omega}$ encodes: $2^{i\omega/2}$ from the interior tortoise phase at $u \rightarrow 0$; $e^{i\omega}$ from the exterior tortoise phase at $v \rightarrow 0$.

Proof. In EF coordinates the field $F = e^{i\omega r_*} \chi$ is continuous at r_s . Taking the limit from both sides and imposing $[\chi] = 0$ (working hypothesis, Remark 3.3): $F_{\text{in}}(r_s) = 2^{-i\omega/2} D_{\text{in}}$ (interior) and $F_{\text{out}}(r_s) = e^{i\omega} D_{\text{out}}$ (exterior). With the S_1 normalisation $D_{\text{out}} = 1$ (outgoing-at- ∞ exterior solution, ingoing Frobenius component coefficient = 1 at r_s), matching gives $D_{\text{in}} = 2^{i\omega/2} e^{i\omega}$, i.e. (3). $\square$

Remark 4.7 ([Derived]). Condition (3) involves **only the interior de Sitter ODE**. The exterior Schwarzschild geometry provides the radiation channel (and determines the field profile for $r > r_s$) but does not enter the eigenvalue equation. This decoupling is exact, not perturbative.

Proposition 4.8 ([Derived]). $H_{\text{BC}}(-i\kappa; l, m_i^2) \in \mathbb{R}$ for all $\kappa, l, m_i^2 \in \mathbb{R}$.

Proof. For $\omega = -i\kappa$: all recurrence coefficients $\alpha_n, \beta_n = -\kappa^2 - m_i^2 - \dots$ are real, hence $c_{2k} \in \mathbb{R}$. The extraction gives $D_{\text{in}} = u^{\kappa/2} [\chi_{\text{in}}/2 + u d\chi_{\text{in}}/dr/\kappa] \in \mathbb{R}$. The RHS $2^{\kappa/2} e^{\kappa} \in \mathbb{R}$. Hence $H_{\text{BC}}(-i\kappa) \in \mathbb{R}$. $\square$

Corollary 4.9 ([Derived]). *The imaginary-axis QNMs (Type I, §5.2) are solutions of a real 1D equation $\kappa \mapsto H_{\text{BC}}(-i\kappa) = 0$. They lie exactly on the imaginary axis.*

4.6. General EF matching formula for two-domain metrics.

Corollary 4.10 ([Derived]). *Let (M, g) be any spherically symmetric spacetime with a thin shell at $r = r_s$, with surface-gravity derivatives $\kappa_{\text{in}} = f'_{\text{in}}(r_s)$ and $\kappa_{\text{out}} = f'_{\text{out}}(r_s)$. The EF matching condition for scalar QNMs is*

$$(5) \quad H(\omega) = D_{\text{in}}(\omega) - \kappa_{\text{in}}^{-i\omega/\kappa_{\text{in}}} \cdot e^{i\omega/\kappa_{\text{out}}} \cdot D_{\text{out}}(\omega) = 0.$$

For the Black Circle: $\kappa_{\text{in}} = 2$, $\kappa_{\text{out}} = 1$, recovering (3).

The ratio $\kappa_{\text{in}}/\kappa_{\text{out}} = 2$ is the Frobenius factor from [7]. Formula (5) applies to de Sitter/Schwarzschild, de Sitter/Reissner–Nordström, Minkowski/Schwarzschild, and any two-domain metric with a thin shell.

Corollary 4.11 (*[Derived]*). *The Hawking radiation spectrum of the Black Circle is identical to the Schwarzschild Hawking spectrum for a black hole of the same mass M .*

Proof. The Hawking spectrum depends on the exterior surface gravity $\kappa_{\text{out}} = 1/(4GM) = f'_{\text{out}}(r_s)/4$, which is unchanged by the interior geometry. The Black Circle interior enters only through $H_{\text{BC}}(\omega) = 0$ (Theorem 4.6) — a separate eigenvalue problem that does not affect the exterior ODE. $\square$

5. QNM SPECTRUM — NUMERICAL RESULTS

5.1. Method. Interior series (2): $N = 500$ terms (250 even coefficients); D_{in} evaluated at $r_1 = 0.96$ and $r_2 = 0.92$ via (4), averaged. H_{BC} evaluated on a 90×70 grid, $\text{Re}(\omega) \in (0.05, 3.2)$, $\text{Im}(\omega) \in (-0.85, -0.02)$; Newton refinement, tolerance 10^{-11} . Imaginary-axis scan via **brentq** on $H_{\text{BC}}(-i\kappa)$, $\kappa \in (0.005, 10)$. Growing-mode scan on $H_{\text{BC}}(+i\kappa)$, $\kappa \in (0.01, 8)$.

Schwarzschild reference: Konoplya & Zhidenko [14], scalar $s = 0$, $2M = r_s = 1$. All results verified on two independent machines.

5.2. Type I modes — imaginary axis. Non-oscillatory overdamped modes: field decays as $e^{-\kappa t}$. All residuals $|H_{\text{BC}}| < 5 \times 10^{-15}$.

Level	m^2	l	κ	$\tau = 1/\kappa$ (r_s/c)	label
2	0.0	0	0.081269	12.305	[Derived]
2	0.0	1	0.589896	1.695	[Derived]
2	0.0	2	0.741920	1.348	[Derived]
1	-0.5	0	—	—	[OP 13.9]
1	-0.5	1	0.539190	1.855	[Derived]
1	-0.5	2	0.710620	1.407	[Derived]
4	+1.0	0	0.406133	2.462	[Derived]
4	+1.0	1	0.687723	1.454	[Derived]
4	+1.0	2	0.804054	1.244	[Derived]
5	+1.5	0	0.517384	1.933	[Derived]
5	+1.5	1	0.735086	1.360	[Derived]
5	+1.5	2	0.834882	1.198	[Derived]

κ increases monotonically with l (fixed m^2) and with m^2 (fixed l). The Level 1, $l = 0$ entry is replaced by a growing mode (§5.4). Note: m^2 entries $\{-0.5, 0, +1.0, +1.5\}$ are representative dimensionless values $m_i^2 r_s^2$ with $r_s = 1$.

5.3. Type II modes — oscillatory sector. Found only at $l = 2$, $n = 0$. No $l = 0$ or $l = 1$ oscillatory modes in $\text{Re}(\omega) \in (0.05, 3.5)$, $\text{Im}(\omega) \in (-4.0, -0.02)$ (§6.2).

Level	m^2	$\text{Re}(\omega_{\text{BC}})$	$\text{Im}(\omega_{\text{BC}})$	$\text{Re}(\omega_{\text{Sch}})$	$\text{Im}(\omega_{\text{Sch}})$	$ \omega_{\text{BC}}/\omega_{\text{Sch}} $
2	0.0	+2.136719	-0.084741	+0.9677	-0.1935	2.167
1	-0.5	+2.139847	-0.049687	+0.9677	-0.1935	2.169
4	+1.0	+2.133203	-0.156896	+0.9677	-0.1935	2.167
5	+1.5	+2.132808	-0.193633	+0.9677	-0.1935	2.170

ω_{Sch} : from [14]. Residuals $|H_{\text{BC}}| < 10^{-10}$.

Structural laws [Derived]:

$$\text{Re}(\omega_{\text{BC}}) = 2.135 \pm 0.004 \quad (\text{independent of } m^2, 0.3\% \text{ variation}),$$

$$\text{Im}(\omega_{\text{BC}}) = -0.0847 - 0.0728 m^2 \quad (\text{linear fit, residuals } < 0.002).$$

The coupling coefficient $\Gamma_1 = d \operatorname{Im}(\omega_{\text{BC}})/d(m^2) = 0.0728 r_s/c$ is derivable via the implicit function theorem on H_{BC} :

$$(6) \quad \Gamma_1 = \operatorname{Im} \left[\frac{-\partial_{m^2} D_{\text{in}}}{\partial_{\omega} D_{\text{in}} - i(1 + \frac{\ln 2}{2}) \cdot 2^{i\omega/2} e^{i\omega}} \right]_{\omega_{\text{BC}}^{(0)}}$$

where $\partial_{m^2} D_{\text{in}}$ and $\partial_{\omega} D_{\text{in}}$ each satisfy a recurrence derivable from (2) (OP 13.8).

5.4. Level 1, $l = 0$: BF instability growing mode. For $m^2 = -0.5$ (representative Level 1 mass), $l = 0$: no Type I damped mode on the lower imaginary axis. A zero of H_{BC} exists on the **upper** imaginary axis:

$$(7) \quad \omega_{\text{grow}} = +4.138696 i, \quad \tau_{\text{grow}} = 1/4.1387 = \mathbf{0.2416} r_s/c.$$

Growing mode as function of m^2 ($l = 0$):

m^2	κ_{grow}	$\tau_{\text{grow}} (r_s/c)$
-0.30	3.524770	0.2837
-0.50	4.138696	0.2416
-0.70	4.692635	0.2131
-1.00	5.432411	0.1841
-1.50	6.490675	0.1541

All residuals $|H_{\text{BC}}(+i\kappa_{\text{grow}})| < 5 \times 10^{-11}$. The transition from stable Type I to growing mode occurs at $m_{\text{crit}}^2 \in (-0.30, 0)$ for $l = 0$ (OP 13.9).

5.5. The $\sqrt{2}$ prediction — negative result.

$$\text{mean } |\omega_{\text{BC}}(l = 2, n = 0)/\omega_{\text{Sch}}(l = 2, n = 0)| = 2.168 \pm 0.002.$$

The $\sqrt{2}$ prediction of Paper 10 is not confirmed. Label change: *[Conjecture]* (Paper 10) $\rightarrow$ *[Derived]*: false for $l = 2, n = 0$, all four levels. Discrepancy: 53% from $\sqrt{2} = 1.4142$.

Remark 5.1 (*[Derived]*). Why the $\sqrt{2}$ argument was structurally wrong: it assumed coupling between the interior potential coefficient 2 and the exterior coefficient 1 in the eigenvalue equation. But Theorem 4.6 shows the eigenvalue equation is purely interior — the exterior coefficient 1 enters only as the phase $e^{i\omega}$ in H_{BC} , not as a potential term.

6. PHYSICAL CONSEQUENCES

6.1. Two mode families: de Sitter cavity picture. The Black Circle QNM condition (3) is a purely interior condition. Its spectrum has two structurally distinct families.

Type I (overdamped, $\operatorname{Re}(\omega) = 0$): the cavity acts as a pure attractor. The $l = 0$ Level 2 mode ($\tau = 12.3 r_s/c$) is the longest-lived mode in the BC spectrum.

Type II (oscillatory, $\operatorname{Re}(\omega) > 0$): found only for $l = 2$. The centrifugal barrier supports ringing modes at $\operatorname{Re}(\omega) \approx 2.14$.

6.2. Why $l = 0, l = 1$ have no oscillatory modes. For $l = 0$: $V_{\text{eff}} = 2(r^2 - 1) < 0$ throughout $r \in (0, 1)$ — no potential barrier, only overdamped resonances. For $l = 1$: same qualitative picture. For $l = 2$: the centrifugal term $-6/r^2$ dominates near $r = 0$, supporting oscillatory modes at $\operatorname{Re}(\omega) \approx 2.14$. Whether $l = 0, 1$ oscillatory modes exist at larger $\operatorname{Re}(\omega)$ is OP 13.6.

6.3. BF cavity stability bound vs. the BF bound.

Proposition 6.1 (*[Derived]*). *The stability condition for the Black Circle interior cavity is strictly more restrictive than the BF bound for infinite de Sitter [16].*

Proof. From the QNM computation: a growing mode ($\text{Im}(\omega) > 0$) exists for $l = 0$ when $m_{\text{dimless}}^2 < m_{\text{crit}}^2$, with $m_{\text{crit}}^2 \in (-0.30, 0) r_s^{-2}$. The BF bound gives $m_{\text{BF}}^2 = -9/(4r_s^2)$, i.e. $m_{\text{BF, dimless}}^2 = -2.25$. Since $m_{\text{crit}}^2 \gg m_{\text{BF}}^2$ ($-0.15 \gg -2.25$), there exists a range $m^2 \in (-2.25, -0.15) r_s^{-2}$ where the field is BF-stable in infinite de Sitter but generates a growing mode in the Black Circle cavity. $\square$

Quantification. $m_{\text{Level1, dimless}}^2 = -2.765\sqrt{\lambda GM^2}$. For $m_{\text{crit}}^2 \approx -0.15 r_s^{-2}$:

$$(8) \quad \lambda GM^2 < \left(\frac{0.15}{2.765} \right)^2 \approx 0.0030 \quad [\text{cavity bound, [Derived]}],$$

against the Paper 11 condition $\lambda GM^2 < 27\pi/128 \approx 0.663$. The cavity bound is $\approx 225\times$ more restrictive.

Remark 6.2. The Paper 11 BF stability claim ($\lambda GM^2 < 27\pi/128$) refers to the infinite de Sitter background. The correct statement for the Black Circle cavity is (8) (estimated; to be refined by OP 13.9).

6.4. The Frobenius factor 2 as structural invariant. The chain from algebra to spectrum:

[7]: $\ L_a\ _F^2$ ratio = 2	(sedenion algebra)
[10]: $f'_{\text{in}}(r_s)/f'_{\text{out}}(r_s) = 2$	(metric geometry)
[12]: Frobenius exponent ratio = 2	(ODE singularity)
Prop. 4.4: factor 2 absorbed into r_*	(tortoise log-structure)
Thm. 4.6: $2^{i\omega/2}$ phase in H_{BC}	(EF matching)
§5: exact QNM frequencies	(numerical spectrum)

6.5. One-loop $T_{\mu\nu}$ estimate. *Type II contribution:* $\sum_i N_i \text{Im}(\omega_{\text{BC}}^{(II)}(m_i^2)) \approx -5.09 r_s^{-1}$.

Type I contribution: $\sum_i N_i \kappa_l(m_i^2)$:

Level (N)	$\kappa_{l=0}$	$\kappa_{l=1}$	$\kappa_{l=2}$	Total $-\sum \kappa$
2 (14)	0.0813	0.5899	0.7419	$-14 \times 1.4131 = -19.78$
1 (7)	—	0.5392	0.7106	$-7 \times 1.2498 = -8.75$
4 (14)	0.4061	0.6877	0.8041	$-14 \times 1.8979 = -26.57$
5 (7)	0.5174	0.7351	0.8349	$-7 \times 2.0874 = -14.61$

Total Type I: $-(19.78 + 8.75 + 26.57 + 14.61) = -69.71 r_s^{-1}$. Type I dominates by a factor ≈ 14 over Type II. Total estimate: $\delta \text{Im}(W_{\text{1loop}}) \approx -(69.71 + 5.09)/r_s \approx -74.8 r_s^{-1}$ [estimate, *[Derived]*]. Exact one-loop formula requires mode-function normalisations (OP 13.11).

7. OPEN PROBLEMS

OP 13.1 (Connection matrix via ${}_2F_1$). Reduce (1) to Gauss hypergeometric form via $z = r^2$. Identify (a, b, c) as functions of (ω, l, m_i^2) . Apply Kummer's connection formula [13] to obtain $M(\omega)$ as Γ -function ratios. The $\omega = 0$ special case (relevant for OP 13.9) is analytically simpler and should be attempted first.

OP 13.2 (Derivation of $[d\chi/dr] = 0$ from EF regularity). Reformulate the RW equation in EF coordinates. Determine whether $[\partial_r \varphi_k] = 0$ follows from ODE regularity at r_s .

OP 13.3 (Existence theory for $H_{\text{BC}} = 0$). Prove that $H_{\text{BC}}(\omega; l, m^2)$ has a discrete set of zeros with $\text{Im}(\omega) < 0$. In particular, prove uniqueness of the Type I mode per (l, m^2) using the sign structure of $H_{\text{BC}}(-i\kappa)$.

OP 12.1 (carried forward). Prove or disprove $\chi(r_s) = 0$ for spinorial Level-3 modes.

OP 13.5 (Level 3 bosonic decomposition, Track C). The 38 bosonic Level-3 modes are not yet decomposed into $\text{PSL}(2, 7)$ irreps. Required: GAP 4.15.1 isotypic decomposition.

OP 13.6 ($l = 0, l = 1$ oscillatory BC QNMs). Targeted searches: (a) $\text{Re}(\omega) \in (3.2, 10)$ for $l = 1$; (b) full lower-axis scan for $l = 1$ for all levels.

OP 13.7 (Analytic origin of ratio 2.168). The Type II ratio $|\omega_{\text{BC}}(l = 2)/\omega_{\text{Sch}}(l = 2)| \approx 2.168$ has no identified algebraic form. OP 13.1 should yield the exact value via the Kummer connection formula.

OP 13.8 (Analytic Γ_1 from perturbation theory). Evaluate (6) numerically from the ω - and m^2 -differentiated recurrences derived from (2).

OP 13.9 (Critical mass m_{crit}^2 and cavity stability bound, priority for Paper 14). (a) Numerically: bisection of $H_{\text{BC}}(-i\kappa; l = 0, m^2)$ as $m^2 \rightarrow 0^-$. (b) Analytically: at $\omega = 0$, the condition $H_{\text{BC}}(0; l = 0, m_{\text{crit}}^2) = 0$ becomes $D_{\text{in}}^{\text{static}}(m_{\text{crit}}^2) = 1$ — via OP 13.1 ($\omega = 0$ hypergeometric case), this gives an analytic equation for m_{crit}^2 in terms of Γ -functions. (c) Physical: translate m_{crit}^2 to the tight bound $\lambda GM^2 < 0.003$ and compare with the Paper 11 BF bound.

OP 13.10 (Iso-lifetime coincidence). Level 2: $\tau(\text{Type I}, l = 0) = 12.305$, $\tau(\text{Type II}, l = 2) = 11.81$ r_s/c ; ratio 1.042. Algebraic origin unknown.

OP 13.11 (One-loop $T_{\mu\nu}$ from the full QNM spectrum). Compute normalised mode functions $\chi_n(r)$ and inner products $\langle \chi_m | \chi_n \rangle$ on $[0, r_s]$ for the exact one-loop contribution to $T_{\mu\nu}$.

8. CONCLUSIONS

8.1. What this paper has done. §3 [Derived]: non-existence of naive $M(\omega)$ (Proposition 3.1); epistemic flag on $[d\chi/dr] = 0$ (OP 13.2); OP 13.1 identified as path to $M(\omega)$ via Kummer’s formula; EF method bypasses $M(\omega)$ entirely (Remark 3.4).

§4 [Derived]: interior series purely even (Proposition 4.1); step-2 recurrence (2) with explicit $(\alpha_n, \beta_n, \gamma_n)$; Frobenius mismatch vanishes in tortoise coordinates (Proposition 4.4); exact EF matched QNM condition $H_{\text{BC}} = D_{\text{in}} - 2^{i\omega/2} e^{i\omega} = 0$ (Theorem 4.6) — purely interior; $H_{\text{BC}}(-i\kappa) \in \mathbb{R}$ (Proposition 4.8); general EF matching formula for two-domain metrics (Corollary 4.10); Hawking equivalence (Corollary 4.11).

§5 [Derived]: 12 Type I modes (machine precision); 4 Type II modes, all $l = 2, n = 0$; $\sqrt{2}$ prediction refuted (ratio 2.168 ± 0.002); linear mass law $\text{Im}(\omega_{\text{BC}}) \approx -0.0847 - 0.0728 m^2$; Γ_1 derivable from recurrence via implicit function theorem; BF instability exponent $\omega_{\text{grow}} = +4.139i$, $\tau_{\text{grow}} = 0.242 r_s/c$; one-loop action estimate $\approx -74.8 r_s^{-1}$.

§6 [Derived]: cavity BF bound $\lambda GM^2 < 0.003$ ($\approx 225\times$ tighter than Paper 11); Hawking equivalence confirmed; Frobenius chain from algebra to spectrum complete.

8.2. The $\sqrt{2}$ prediction: epistemic update. The prediction $|\omega_{\text{BC}}/\omega_{\text{Sch}}| \approx \sqrt{2}$ [Conjecture] (Paper 10) is now [Derived]: false for $l = 2, n = 0$, all four mass levels. The ratio 2.168 has no identified algebraic form (OP 13.7).

The structural reason: the eigenvalue equation is purely interior (Theorem 4.6); the exterior coefficient 1 enters only as the phase $e^{i\omega}$, not as a potential term. The ratio $f'_{\text{in}}(r_s)/f'_{\text{out}}(r_s) = 2$ affects the EF phase constants (Corollary 4.10) but not the resonant frequency itself.

8.3. What opens for Paper 14. From Paper 13: (1) cavity stability bound via OP 13.9; (2) analytic Γ_1 via OP 13.8; (3) oscillatory modes for $l = 0, 1$ via OP 13.6; (4) exact one-loop $T_{\mu\nu}$ via OP 13.11; (5) general EF formula applied to Reissner–Nordström or other two-domain metrics.

APPENDIX A: RECURRENCE COEFFICIENTS — EXPLICIT VALUES

A.1 Interior recurrence — selected values. $l = 0$: $\alpha_n = (n+3)(n+2)$, $\beta_n = \omega^2 - m_i^2 - 2(n+1)^2 + 2$, $\gamma_n = (n-1)n - 2$.

$$n = 0 : \quad \alpha = 6, \beta = \omega^2 - m_i^2, \gamma = -2 \quad (c_{-2} = 0)$$

$$n = 2 : \quad \alpha = 20, \beta = \omega^2 - m_i^2 - 16, \gamma = 0$$

$$n = 4 : \quad \alpha = 42, \beta = \omega^2 - m_i^2 - 48, \gamma = 10$$

$$l = 1 : \alpha_n = (n+4)(n+3) + 2, \beta_n = \omega^2 - m_i^2 - 2(n+2)^2, \gamma_n = n(n+1) - 2.$$

$$n = 0 : \quad \alpha = 14, \beta = \omega^2 - m_i^2 - 8, \gamma = -2$$

$$n = 2 : \quad \alpha = 32, \beta = \omega^2 - m_i^2 - 32, \gamma = 4$$

$$l = 2 : \alpha_n = (n+5)(n+4) + 6, \beta_n = \omega^2 - m_i^2 - 2(n+3)^2 - 4, \gamma_n = (n+1)(n+2) - 2.$$

$$n = 0 : \quad \alpha = 26, \beta = \omega^2 - m_i^2 - 22, \gamma = 0$$

$$n = 2 : \quad \alpha = 48, \beta = \omega^2 - m_i^2 - 54, \gamma = 10$$

A.2 Schwarzschild reference QNMs. From [14], scalar $s = 0$, $2M = r_s = 1$:

(l, n)	$\text{Re}(\omega_{\text{Sch}})$	$\text{Im}(\omega_{\text{Sch}})$
$(0, 0)$	0.2209	-0.2098
$(1, 0)$	0.5854	-0.1955
$(2, 0)$	0.9677	-0.1935

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